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$\Lambda$  AND  $K_s^0$  PRODUCTION IN  $p+C$   
COLLISIONS AT 10 GeV/ $c$

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Экспериментальные данные, полученные с помощью 2-м пропановой пузырьковой камеры, проанализированы для канала реакции  $pC \rightarrow \Lambda(K_s^0)X$  при импульсе пучка протонов 10 ГэВ/ $c$ . Экспериментальные измерения инклузивных сечений для выходов  $\Lambda$  и  $K_s^0$  в столкновениях  $p^{12}C$  равны  $\sigma_\Lambda = (13,3 \pm 1,7)$  мб и  $\sigma_{K_s^0} = (3,8 \pm 0,6)$  мб соответственно.

Также измерены отношения для средних множественностей  $\Lambda/\pi^+$  и  $K_s^0/\pi^+$  для реакции  $pC$ . Экспериментально получаем, что величина для отношений  $\Lambda/\pi^+((5,3 \pm 0,8) \cdot 10^{-2})$  в реакции  $pC$  приблизительно в два раза больше, чем это же отношение, рассчитанное моделью FRITIOF при одинаковых экспериментальных условиях. Это отклонение для отношений средних множественностей увеличивается и достигает максимальных значений ( $\approx$  в 4 раза больше) при импульсах пучка 10–15 ГэВ/ $c$  не только в столкновениях тяжелых ионов, но и в реакции  $C+C$ , как и предсказывает термальная статистическая адронная модель.

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The experimental data from the 2-m propane bubble chamber have been analyzed for  $pC \rightarrow \Lambda(K_s^0)X$  reactions at 10 GeV/ $c$ . The estimation of experimental inclusive cross sections for  $\Lambda$  and  $K_s^0$  production in the  $p^{12}C$  collision is equal to  $\sigma_\Lambda = (13.3 \pm 1.7)$  mb and  $\sigma_{K_s^0} = (3.8 \pm 0.6)$  mb, respectively.

The measured  $\Lambda/\pi^+$  ratio from  $pC$  reaction is equal to  $(5.3 \pm 0.8) \cdot 10^{-2}$ . The experimental  $\Lambda/\pi^+$  ratio in the  $pC$  reaction is approximately two times larger than the  $\Lambda/\pi^+$  ratio received with FRITIOF model for the same reaction and energy. The  $\Lambda/\pi^+$  ratio has significant enhancement for heavy-ion and also for  $C+C$  collisions at 10–15 ГэВ/ $c$  as a thermal statistical hadron model predicted.

The investigation has been performed at the Veksler and Baldin Laboratory of High Energies, JINR.

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## 1. INTRODUCTION

Strangeness enhancement has been extensively discussed as a possible signature for the quark–gluon plasma (QGP) [1, 2]. Strange-particle production has also been analyzed regarding such reaction mechanisms as the multinucleon effect [3], the fireball effect [4], and the deconfinement signal within the context of thermal equilibration models [5–8] (Fig. 3).

In particular, strange particles have been observed extensively in hadron–nucleus and nucleus–nucleus collisions in 4–15-Gev regions [9–14]. The strange-hyperon yields [9–11] are therefore of great interest as an indicator of strange-quark production. The number of  $\Lambda$ 's produced in  $\bar{p}$ +Ta reaction at 4 GeV/c was 11.3 times larger than that expected from the geometrical cross section [9]. In experiments with Si+Au and Au+Au collisions at 11.6 [13] and 14.6 A GeV/c [14] a  $K^+/\pi^+$  ratio in heavy-ion reactions was measured to be 4–5 times larger than the  $K^+/\pi^+$  ratio in  $p + p$  reactions at the same energy. The thermal model [6] gives a good description of  $K^+/\pi^+$ ,  $\Lambda/\pi^+$  ratios for data Au+Au, Si+Au interactions at momentum 10–15 A Gev/c, showing a broad maximum at the same energies.

However, there have not been sufficient experimental data concerning strange-hyperon production over 10–40-GeV/c momentum range. In this paper, new results are presented — the measured inclusive cross sections for  $\Lambda K_s^0$  production and  $\Lambda/\pi^+$  ratio in the reaction  $p+^{12}C$ .

## 2. EXPERIMENTAL PROCEDURE

**2.1. Method.** Experimental data on  $\approx 700\,000$  stereo photographs by the JINR 2-m propane bubble chamber exposed to a 10-GeV/c proton beam [15–21] were analyzed. The primary proton beams must satisfy conditions:  $|tg\alpha| < 0.02$ ,  $1.62 < \beta < 1.69$  rad. The magnetic field ( $B=15.2$  kG) measurement error is  $\Delta B/B = 1\%$ . The fit GRIND-based program GEOFIT [18, 15] is used to measure the kinematics track parameters  $p$ ,  $\alpha$ ,  $\beta$ . Measurements were repeated three times for events which failed in reconstruction by GEOFIT [15].

The estimation of ionization for charged tracks and length for stopped particles permitted one to identify them over the following momentum ranges: protons of  $0.150 \leq p \leq 0.900$  GeV/c and  $K^\pm$  of  $p \leq 0.6$  GeV/c.

**2.2. Identification of  $\Lambda$  and  $K_s^0$ .** The events with  $V^0$  ( $\Lambda$  and  $K_s^0$ ) were identified using the following criteria [19, 20, 21]: 1)  $V^0$  stars from the photographs were selected according to  $\Lambda \rightarrow \pi^- + p$ , neutral  $K_s \rightarrow \pi^- + \pi^+$  or  $\gamma \rightarrow e^+ + e^-$  hypothesis. A momentum limit of  $K_s^0$  and  $\Lambda$  is greater than 0.1 and 0.2 GeV/c, respectively; 2)  $V^0$  stars should have the effective mass of  $K_s^0$  and  $\Lambda$ ; 3) these  $V^0$  stars are directed to some vertices (coplanarity); 4) they should have one vertex, a three-constraint fit for the  $M_K$  or  $M_\Lambda$  hypothesis and after the fit  $\chi_{V^0}^2$  should be selected over the range less than 12; 5) the analysis has shown [20] that the events with undivided  $\Lambda K_s^0$  were assumed to be  $\Lambda$ .

Table 1 presents the number of experimental  $V^0$  events (70% of all events have been identified) produced from interactions of: a) primary beam protons, b) secondary charged particles, and c) secondary neutral particles.

**Table 1. The number of  $V^0$  events from interactions of different types which were registered in stereo photographs without restrictions over effective ranges of propane bubble chamber**

Channel	The number of events from interactions			Total events
	Primary beam protons	Sec. charged particles	Sec. neutral particles	
$\rightarrow \Lambda(\text{only})x$	5276	2814	1063	9387
$\rightarrow K_s^0(\text{only})x$	4122	1795	481	6543
$\rightarrow (\Lambda \text{ and } K_s^0)x$	3381	1095	376	4608

The  $V^0$ 's can be classified into three grades. The first grade comprised  $V^0$ 's which could be identified with the above criteria (1–4) and bubble densities of the positive track emitted from  $V^0$ 's. The second grade comprised  $V^0$ 's which could be undivided  $\Lambda K_s^0$ . For correct identification of the undivided  $V^0$ 's, the  $\alpha$  Armenteros parameter and the  $\cos\theta_{\pi^-}^*$  distributions in the rest frame  $V^0$  (Fig. 1) are used:

$$\alpha = (P_{||}^+ - P_{||}^-)/(P_{||}^+ + P_{||}^-),$$

where  $P_{||}^+$  and  $P_{||}^-$  are the momentum components of positive and negative charged tracks from the  $V^0$  relative direction of the  $V^0$  momentum. The  $\theta_{\pi^-}^*$  is the angle between  $\pi^-$  (from  $K_s^0$  decay) and  $V^0$  in  $V^0$  rest frame. The  $\alpha$  (Fig. 1, a) and  $\cos\theta_{\pi^-}^*$  distributions from  $K_s^0$  decay were isotropic in the  $K_s^0$  rest frame after removing undivided  $\Lambda K_s^0$ . Then these  $\Lambda K_s^0$  events are assumed to be  $\Lambda$  events. In Fig. 1, c we show that the  $\cos\theta_{\pi^-}^*$  distributions for the  $\Lambda + \Lambda K_s^0$  have been also isotropic in  $V^0$  rest frame. As a result of the above procedure, 8.5% of  $K_s^0$  are lost and 4.6% of  $K_s^0$  are interpreted as  $\Lambda$ . The third grade comprised  $V^0$ 's which could be the invisible  $V^0$ 's at a large azimuth angle  $\phi$  [20]. The average  $\phi$  weights were  $\langle w_\phi \rangle = 1.06 \pm 0.02$  for  $K_s^0$  and  $\langle w_\phi \rangle = 1.14 \pm 0.02$  for  $\Lambda$ .

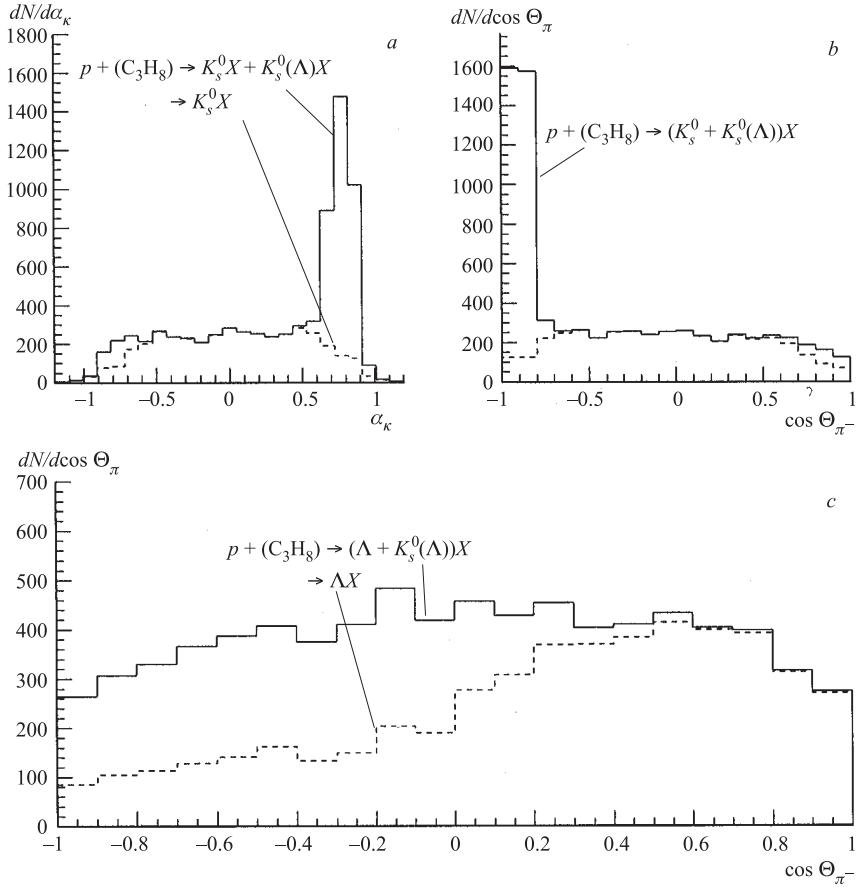


Fig. 1. Distributions of  $\alpha$  (Armenteros parameter) and  $\cos \Theta^*$  are used for correct identification of the undivided  $V^0$ 's.  $\alpha = (P_\parallel^+ - P_\parallel^-)/((P_\parallel^+ + P_\parallel^-)$ , where  $P_\parallel^+$  and  $P_\parallel^-$  are the parallel components of momenta of positive and negative charged tracks with respect to the direction of momentum  $K_s^0(\Lambda)$ ,  $\cos \Theta^*$  is the angular distribution of  $\pi^-$  from  $K_s^0$  decay. Distributions of  $\alpha$  and  $\cos \Theta$  were isotropic in the rest frame of  $K_s^0$  when undivided  $K_s^0(\Lambda)$  were assumed to be  $\Lambda$

Figures 2, a, c and 2, b, d show the effective mass distribution of  $\Lambda$  (8657 events),  $K^0$  (4122 events) particles and their  $\chi^2$  from kinematics fits, respectively, produced from the beam protons interacting with propane targets. The measured masses of these events have the following Gaussian distribution parameters  $\langle M(K_s) \rangle = 497.7 \pm 3.6$ , SD = 23.9 MeV/ $c^2$  and  $\langle M(\Lambda) \rangle = 1117.0 \pm 0.6$ , SD = 10.0 MeV/ $c^2$ . The masses of the observed  $\Lambda$ ,  $K_s^0$  are consistent with their PDG values. The expected functional form for  $\chi^2$  is depicted with the dotted histogram (Fig. 2).

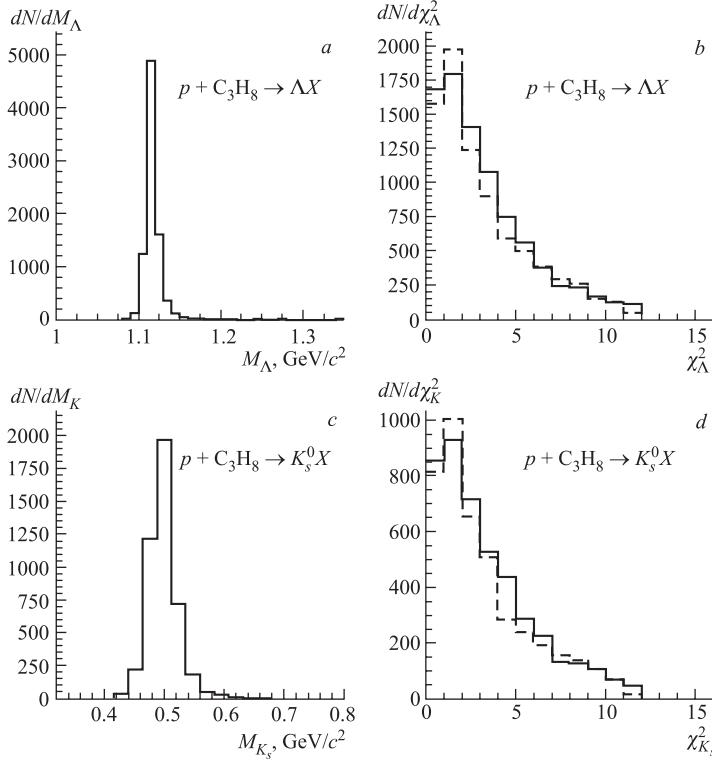


Fig. 2. The distribution of experimental  $V^0$  events produced from interactions of beam protons with propane: *a*) for the effective mass of  $M_\Lambda$ ; *b*) for  $\chi^2_\Lambda$  of the fits via the decay mode  $\Lambda \rightarrow \pi^- + p$ ; *c*) for the effective mass of  $M_{K_s^0}$ ; *d*) for  $\chi^2_{K_s^0}$  of the fits via decay mode  $K_s^0 \rightarrow \pi^- + \pi^+$ . The expected functional form for  $\chi^2$  is depicted with the dotted histogram

Each  $V^0$  event weighted by a factor  $w_{\text{geom}} = 1/e_\tau$ , where  $e_\tau$  is the probability for potentially observed  $V^0$ , can be expressed as

$$e_\tau = \exp(-L_{\min}/L) - \exp(-L_{\max}/L), \quad (1)$$

where  $L = cp\tau/M$  is the flight length of the  $V^0$ ,  $L_{\max}$  is the path length from the reaction point to the boundary of fiducial volume, and  $L_{\min}$  (0.5 cm) is an observable minimum distance between the reaction point and the  $V^0$  vertex.  $M$ ,  $\tau$ , and  $p$  are the mass, lifetime, and momentum of the  $V^0$ . The average geometrical weights are  $1.34 \pm 0.02$  for  $\Lambda$  and  $1.22 \pm 0.04$  for  $K^0$ .

Now, let us examine a possibility from neutron stars of imitating  $\Lambda$  and  $K_s^0$  by using FRITIOF model [22] for the hypothesis reaction  $p+C \rightarrow n + X$ ,

$n + n \rightarrow \pi^- p(\pi^-\pi^+) + X^0$  including Fermi motion in carbon. Then, these background events were analyzed by using the same experimental conditions for the  $V^0$  selection. The 2-vertex analysis has shown that the background from neutron stars is equal to 0.1 % for  $\Lambda$  and 0.001 for  $K_s^0$  events.

**2.3. The Selection of Interactions on Carbon Nucleus.** The criteria for selection of interaction with carbon has been shown [19,25]. The  $p+C \rightarrow \Lambda(K_s^0)X$  reaction was selected by the following criteria:

1.  $Q = n_+ - n_- > 2$ ;
2.  $n_p + n_\Lambda > 1$ ;
3.  $n_p^b + n_\Lambda^b > 0$ ;
4.  $n_- > 2$ ;
5.  $n_{\text{ch}} = \text{odd number}$  ;
6.  $\frac{E_p(\Lambda) - P_{p(\Lambda)} \cos \Theta_{p(\Lambda)}}{m_t} > 1$ ;

$n_+$  and  $n_- >$  are the numbers of positive and negative particles on the star;  $n_p$  and  $n_\Lambda$  are the numbers of protons and  $\Lambda$  hyperons with momentum  $P < 0.75$  GeV/c on the star;  $n_p^b$  and  $n_\Lambda^b$  are the numbers of protons and  $\Lambda$  hyperons emitted in backward direction;  $E_p(\Lambda)$ ,  $P_{p(\Lambda)}$  and  $\Theta_{p(\Lambda)}$  are an energy, a momentum and an emitted angle of protons (or  $\Lambda$ 's) in the lab system;  $m_t$  is a mass of the target. These criteria were separated  $\approx 83\%$  from all inelastic  $p+C$  interactions [25]. The  $p+C$  events were selected by the above criteria from simulated  $p+\text{propane}$  interactions by using FRITIOF model [22] under experimental conditions and by the above criteria. Results of the simulation show loss of 18% and 20% from interactions  $pC \rightarrow \Lambda X$  and  $pC \rightarrow K_s^0 X$ , respectively. The simulation by FRITIOF model also shows that the contribution from  $pp \rightarrow \Lambda X$  and  $pp \rightarrow K_s^0 X$  in  $pC$  interactions is equal to 1.0% and 0.3%, respectively.

### 3. THE MEASURED CROSS SECTIONS $\Lambda$ AND $K^0$

The cross section is defined by the formula:

$$\sigma = \frac{\sigma_0 N_r^{V^0}}{e} \prod_i w_i = \frac{\sigma_r N_r^{V^0} w_{\text{hyp}} w_{\text{geom}} w\phi w_{\text{kin}} w_{\text{int}}}{N_{\text{int}}^r e_1 e_2 e_3}, \quad (2)$$

where  $e_1$  is the efficiency of search for  $V^0$  on the photographs,  $e_2$  is the efficiency of measurements,  $e_3$  is the probability of decay via the channel of charged particles ( $\Lambda \rightarrow p\pi^-$ ,  $K^0 \rightarrow \pi^+\pi^-$ ),  $\sigma_0 = \sigma_r/N_r$  is the total cross section, where  $\sigma_r$  is the total cross section for registered events,  $N_r$  is the total number of registered interactions of beam protons over the range of the chamber.  $\sigma_t(p + C_3H_8) = 3\sigma_{pC} + 8\sigma_{pp} = (1456 \pm 88)$  mb [27], where  $\sigma_t$ ,  $\sigma_{pC}$  and  $\sigma_{pp}$  are the total cross sections in interactions  $p+C_3H_8$ ,  $p+C$  and  $p+p$ , respectively. The propane bubble

chamber method has allowed the registration of the part of all elastic interactions with the propane [23, 24], therefore the total cross section of registered events is equal to  $\sigma_r(p + C_3H_8) = 3\sigma_{pC}(\text{inelastic}) + 8\sigma_{pp}(\text{inelastic}) + 8\sigma_{pp}(\text{elastic})0.70 = (1049 \pm 60) \text{ mb}$  [23].  $w_i$  are weights for the lost events with  $V^0$  for (Table 2):  $w_{\text{geom}}$  — the  $V^0$  decay outside the chamber;  $w_\phi$  — the required isotropy for  $V^0$  in the azimuthal (XZ) plane;  $w_{\text{hyp}}$  — the undivided  $\Lambda K_s^0$  events;  $w_{\text{int}}$  — the selected as  $p + {}^{12}\text{C}$  from the interaction of  $p + C_3H_8$ ;  $w_{\text{kin}}$  — the kinematic conditions (by FRITIOF);  $w_{\text{int}}$  — the  $V^0 + \text{propane}$  interactions.

**Table 2. Weight of the lost experimental events with  $\Lambda$  and  $K_s^0$  for  $pC$  and  $pp$  interactions**

Type of reaction	$1/e_1$	$1/e_2$	$w_{\text{geom}}$	$w_\phi$	$w_{\text{int}}$	$w_{\text{kin}}$	$1/e_3$	$W_{\text{sum}}$
$pC \rightarrow \Lambda X$	1.14	1.25	1.34	1.14	1.11	1.18	1.56	$4.37 \pm 0.37$
$pp \rightarrow \Lambda X$	1.14	1.25	1.36	1.14	1.11	1.37	1.56	$5.15 \pm 0.44$
$pC \rightarrow K_s^0 X$	1.14	1.25	1.22	1.06	1.04	1.04	1.47	$2.93 \pm 0.25$
$pp \rightarrow K_s^0 X$	1.14	1.25	1.36	1.06	1.05	1.06	1.47	$3.31 \pm 0.28$

Table 3 shows that the experimental cross sections are calculated by (2) for inclusive  $\Lambda$  hyperon and  $K_s^0$  meson productions in the interactions of  $pp$  and  $pC$  at beam momentum of 10 GeV/c.

**Table 3. Cross sections of  $\Lambda$  hyperons and  $K_s^0$  mesons for  $pp$  and  $pC$  interactions at beam momentum of 10 GeV/c**

Type of reaction	$N_{V^0}^{\text{exp}}$	$W_{\text{sum}}$	$N_{V^0}^t$ (total)	$n_{V^0} = N_{V^0}^t/N_{\text{in}}$	$\sigma, \text{mb}$
$pC \rightarrow \Lambda X$	6126	$4.37 \pm 0.37$	26770	$0.053 \pm 0.005$	$13.3 \pm 1.6$
$pp \rightarrow \Lambda X$	836	$5.15 \pm 0.44$	4303	$0.026 \pm 0.003$	$0.80 \pm 0.08$
$pC \rightarrow K_s^0 X$	3188	$2.93 \pm 0.25$	9341	$0.018 \pm 0.002$	$3.8 \pm 0.6$
$pp \rightarrow K_s^0 X$	699	$3.31 \pm 0.28$	2313	$0.015 \pm 0.001$	$0.43 \pm 0.04$

Ratios of average multiplicities of  $\Lambda$  hyperons and  $K_s^0$  mesons to multiplicities of  $\pi^+$  mesons in  $p+C$  interaction at beam momenta 4.2 and 10 GeV/c are shown in Table 4. Experimental data on multiplicity of  $\pi^+$  mesons in the interaction of  $pC$  at momenta 4.2 GeV/c ( $< n_{\pi^+} > = 0.71 \pm 0.01$ ) and 10 GeV/c ( $< n_{\pi^+} > = 1.0 \pm 0.05$ ) are taken from publications [26] and [25], respectively.

The  $\Lambda/\pi^+$  ratio for  $C+C$  reaction is shown in Table 5 and in Fig. 3. This  $\Lambda/\pi^+$  ratio at momentum 10 GeV/c has been obtained by using the Glauber approach on the experimental cross section for  $p+C \rightarrow \Lambda X$  reaction. As can be seen from experimental data in Table 5 and thermal statistical model (Fig. 3)

**Table 4. Ratios of average multiplicities of  $\Lambda$  hyperons and  $K_s^0$  mesons to multiplicities of  $\pi^+$  mesons for  $p+C$  interaction at beam momenta of 4.2 and 10  $\text{GeV}/c$**

	$pC$ experiment (10 $\text{GeV}/c$ )	$pC$ FRITIOF (10 $\text{GeV}/c$ )	$Cp$ experiment (4.2 $\text{GeV}/c$ )	$Cp$ FRITIOF (4.2 $\text{GeV}/c$ )
$\langle n_\Lambda \rangle / \langle n_{\pi^+} \rangle \cdot 10^2$	$5.3 \pm 0.8$	2.6	$0.7 \pm 0.3$	0.9
$\langle n_{K_s^0} \rangle / \langle n_{\pi^+} \rangle \cdot 10^2$	$1.8 \pm 0.3$	1.8	$0.3 \pm 0.2$	0.3

there is a very clearly pronounced enhancement specially in the  $\Lambda/\pi^+$  ratio for hadron–nucleus and nuclear collisions at 10–15  $A$   $\text{GeV}/c$ .

**Table 5. Ratios of average multiplicity of  $\Lambda$  hyperons to multiplicity of  $\pi^+$  mesons for  $C+C$  interactions at beam momenta of 4.2 and 10  $\text{GeV}/c$**

	Experiment (4.2 $\text{GeV}/c$ )	Experiment (10 $\text{GeV}/c$ )
$\langle n_\Lambda \rangle / \langle n_{\pi^+} \rangle \cdot 10^2$	$2.0 \pm 0.6$	$10.9 \pm 1.7$

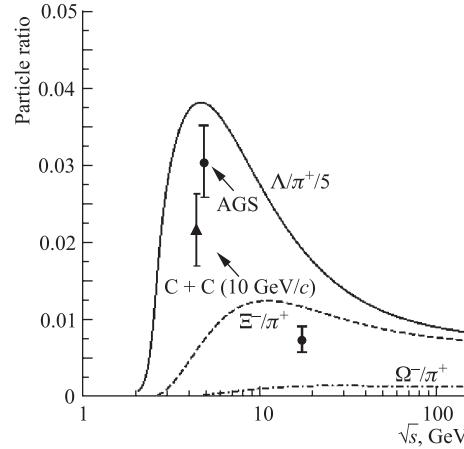


Fig. 3. Prediction of the statistical thermal model [6] for  $\Lambda/\pi^+$  (note the factor 5),  $\Xi^-/\pi^+$  and  $\Omega^-/\pi^+$  ratios versus  $\sqrt{s}$ . For compilation of AGS data see [7]. The  $\Lambda/\pi^+$  ratio in interaction C+C in figure is obtained by using data from this experiment

#### 4. CONCLUSION

The experimental data from the 2-m propane bubble chamber have been analyzed for  $pC \rightarrow \Lambda(K_s^0)X$  reaction at 10 GeV/c. The estimation of experimental inclusive cross sections for  $\Lambda$  and  $K_s^0$  production in  $pC$  collisions is equal to  $\sigma_\Lambda = (13.3 \pm 1.7)$  mb and  $\sigma_{K_s^0} = (3.8 \pm 0.6)$  mb, respectively. The measured  $\Lambda/\pi^+$  ratio in  $pC$  and  $pp$  reactions is equal to  $(5.3 \pm 0.8) \cdot 10^{-2}$  and  $(2.6 \pm 0.4) \cdot 10^{-2}$ , respectively. The experimental  $\Lambda/\pi^+$  ratio in the  $pC$  reaction is approximately two times larger than the  $\Lambda/\pi^+$  ratio received from  $pp$  or  $pC$  reaction by FRITIOF model for the same energy. There is a very clearly pronounced enhancement in the  $\Lambda/\pi^+$  ratio for Au+Au and C+C collisions at 10–15 A GeV/c as predicted by the thermal statistical hadron model.

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